## Master NPAC

Cosmology homework #2

## December 1, 2022

Answers are to be sent by email to nicolas.regnault-AT-lpnhe.in2p3.fr by Thursday December 8th, 2022. Please use the following format (given as an example):

<your name>

1. A 2. B 3. C 4. D 5. A 6. B 7. A

We have seen that, in a homogeneous, isotropic Universe in expansion, the metric, in comoving coordinates  $(t, \chi, \theta, \phi)$  can be written as:

$$ds^{2} = -c^{2}d\tau^{2} = -c^{2}dt^{2} + a^{2}(t)\left(d\chi^{2} + r^{2}(\chi)\left(d\theta^{2} + \sin\theta d\phi^{2}\right)\right)$$
(1)

General Relativity tells us that free particles follow geodesics  $(x^{\mu}(\lambda))$ , which are solution of the geodesics equation:

$$\frac{d^2 x^{\mu}}{d\tau^2} + \Gamma^{\mu}_{\alpha\beta} \frac{dx^{\alpha}}{d\tau} \frac{dx^{\beta}}{d\tau} = 0$$
<sup>(2)</sup>

where the  $\Gamma$  coefficients are the christoffel connections. For the metric above, we remind you that:

$$\Gamma_{00}^{0} = \Gamma_{00}^{i} = \Gamma_{0i}^{0} = 0 \quad \Gamma_{0j}^{i} = \frac{\dot{a}}{a} \delta_{j}^{i} \quad \Gamma_{ij}^{0} = a \dot{a} \tilde{g}_{ij} \quad \Gamma_{11}^{1} = \Gamma_{\chi\chi}^{\chi} = 0 \tag{3}$$

In this homework, we study the motion of a spacecraft that leaves our Galaxy ( $\chi = 0$ ) today ( $t = t_0$ ), with an initial velocity  $\frac{d\ell_{\text{phys}}}{dt} = \frac{da(t)\chi(t)}{dt} = v_0$ . We assume that all its fuel is used to generate the spacecraft's initial momentum, and that it follows a purely radial trajectory: ( $t(\tau), \chi(\tau), \theta = \text{const}, \phi = \text{const}$ ).

**1.** Show that

$$\dot{\chi}(t_0) = \frac{d\chi}{dt}(t_0) = \frac{v_0}{a_0} \tag{4}$$

Use  $-c^2 d\tau^2 = g_{\mu\nu} dx^{\mu} dx^{\nu}$  to express  $d\tau/dt$  as a function of a and  $\chi$ :

- (A)  $\frac{d\tau}{dt} = (1 \frac{a^2 \dot{\chi}^2}{c^2})^{1/2}$ (B)  $\frac{d\tau}{dt} = (1 - a^2 \chi^2)^{1/2}$ (C)  $\frac{d\tau}{dt} = (1 - a\chi/c)^{1/2}$
- (D) None of the above
- **2.** Express the initial values (at  $t_0$ ) of  $dt/d\tau$  and  $d\chi/d\tau$

(A) 
$$\frac{dt}{d\tau}_{|t_0} = 1$$
  $\frac{d\chi}{d\tau} = v_0/a_0$ 

- (B)  $\frac{dt}{d\tau}_{|t_0} = (1 \frac{v_0^2}{c^2})^{-1/2}$   $\frac{d\chi}{d\tau} = \frac{v_0}{a_0} \frac{1}{\sqrt{1 v_0^2/c^2}}$
- (C)  $\frac{dt}{d\tau}_{|t_0} = (1 \frac{v_0^2}{c^2})^{-1/2} \quad \frac{d\chi}{d\tau} = \frac{v_0}{a_0}\sqrt{1 v_0^2/c^2}$
- (D) None of the above

**3.** Derive from the geodesic equation obeyed by the spacecraft trajectory the differential equations followed by  $t(\tau)$  and  $\chi(\tau)$ :

(A)  $\frac{d^2t}{d\tau^2} = -a\dot{a}\left(\frac{d\chi}{d\tau}\right)^2 \quad \frac{d}{d\tau}(a^2\frac{d\chi}{d\tau}) = 0$  $\frac{d^2t}{d\tau} = -a\dot{a}\left(\frac{d\chi}{d\tau}\right)^2 \quad \frac{d}{d\tau}(a^2\frac{d\chi}{d\tau}) = 0$ 

(B) 
$$\frac{d\tau}{d\tau^2} = 1$$
  $\frac{d}{d\tau} (\frac{du}{d\tau} \frac{d\chi}{d\tau}) = 0$ 

(C) 
$$\frac{d^2t}{d\tau^2} = -a\dot{a}\dot{\chi}^2 \qquad \frac{d}{d\tau}(a^2\frac{d\chi}{d\tau}) = 0$$

(D) None of the above

**4.** Derive from the above one single differential equation  $d\chi/dt = f(a, \chi, t)$ . Write the limits of this equation for a slowly moving spacecraft  $(v_0 \to 0)$  and a spacecraft moving at relativistic speeds:  $v_0 \to c$ .

(A) 
$$(v_0 \to 0) \dot{\chi} = \frac{a_0 v_0}{a^2(t)} \quad (v_0 \to c) \dot{\chi} = \frac{c}{a(t)}$$

- (B)  $(v_0 \to 0) \dot{\chi} = \frac{v_0}{a(t)}$   $(v_0 \to c) \dot{\chi} = \frac{c}{a(t)}$ (C)  $(v_0 \to 0) \dot{\chi} = v_0 \dot{a}$   $(v_0 \to c) \dot{\chi} = \frac{a_0}{a^2(t)}$
- (D) None of the above

**5.** We assume that we live in a matter-dominated Universe:  $a(t) = a_0(t/t_0)^{2/3}$ , with  $t_0 = (2/3)H_0^{-1}$ . The spacecraft reaches a galaxy at  $t = t_1$ . What is the comoving distance  $a_0\chi_1$  to this galaxy in the two cases above  $(v_0 \to 0 \text{ and } v_0 \to c)$ ?

(A) 
$$(v_0 \to 0) \ a_0 \chi_1 = \frac{2}{H_0} \left( 1 - \left(\frac{t_0}{t_1}\right)^{2/3} \right) \quad (v_0 \to c) \ a_0 \chi_1 = \frac{2c}{H_0} \left( \left(\frac{t_0}{t_1}\right)^{2/3} - 1 \right)$$
  
(B)  $(v_0 \to 0) \ a_0 \chi_1 = \frac{2a_0}{H_0} \left( 1 - \left(\frac{t_0}{t_1}\right)^{1/3} \right) \quad (v_0 \to c) \ a_0 \chi_1 = \frac{2a_0}{H_0} \left( \left(\frac{t_1}{t_0}\right)^{2/3} - 1 \right)$   
(C)  $(v_0 \to 0) \ a_0 \chi_1 = \frac{2v_0}{H_0} \left( 1 - \left(\frac{t_0}{t_1}\right)^{1/3} \right) \quad (v_0 \to c) \ a_0 \chi_1 = \frac{2c}{H_0} \left( \left(\frac{t_1}{t_0}\right)^{1/3} - 1 \right)$ 

(D) None of the above

6. Same question for a vacuum-energy-dominated Universe:  $a(t) = a_0 \exp H_0(t - t_0)$ (A)  $(v_0 \to 0) \ a_0 \chi_1 = \frac{v_0}{2H_0} \left(1 - e^{-2H_0(t_1 - t_0)}\right) \quad (v_0 \to c) \ a_0 \chi_1 = \frac{c}{H_0} \left(1 - e^{-H_0(t_1 - t_0)}\right)$ 

(B) 
$$(v_0 \to 0) \ a_0 \chi_1 = \frac{v_0}{H_0} \left( 1 - e^{-H_0(t_1 - t_0)} \right) \quad (v_0 \to c) \ a_0 \chi_1 = \frac{c}{2H_0} \left( 1 - e^{-2H_0(t_1 - t_0)} \right)$$

(C) 
$$(v_0 \to 0) \ a_0 \chi_1 = \frac{c}{H_0} \left( 1 - e^{-2H_0(t_1 - t_0)} \right) \quad (v_0 \to c) \ a_0 \chi_1 = \frac{v_0}{H_0} \left( 1 - e^{-H_0(t_1 - t_0)} \right)$$

(D) None of the above

**7.** For  $t \to \infty$  and  $v_0 \to c$  is there a limit to the distance we can travel in (1) a matter-dominated Universe and (2) a vacuum-dominated Universe ?

- (A) Yes Yes
- (B) Yes No
- (C) No Yes
- (D) No No

**8.** With the expansion of the Universe, the energy density of non relativistic matter falls like  $a^{-3}$  and the energy density of radiation falls like  $a^{-4}$ . Assuming  $\Omega_{m,0} = 0.3$  (you know  $\Omega_{r,0}$  already, right ?), plot the evolution of  $\rho_m$  and  $\rho_r$  as a function of z. Draw on the plot the value of  $\rho_m$  at z = 3. At which redshift were  $\rho_m$  and  $\rho_r$  equal ? What was their value ?